

Handwritten Notes
On
X-RAYS



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X-RAYS

Like photon.

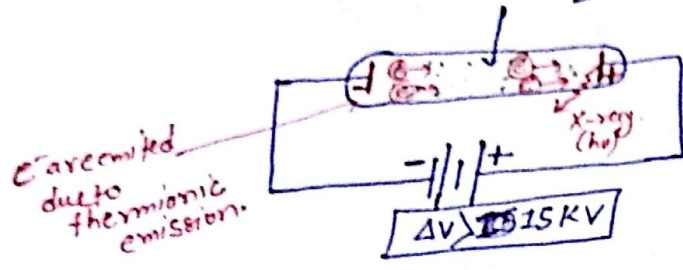
is not deflected from electric & magnetic field.

* Reverse of P.E.E.

condn → * When highly energetic e^s are made to strike metal target, electromagnetic radiation of the order 0.01 Å to 100 Å was observed & known as X-ray.

|| → Röntgen exp: → In presence of high potential & low potential, sum invisible radiation is emitted from anode which is called X-ray.

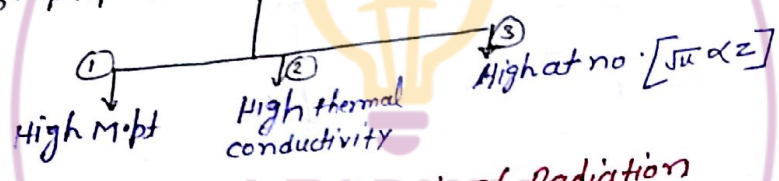
E.D.T [20³ mm of Hg column]



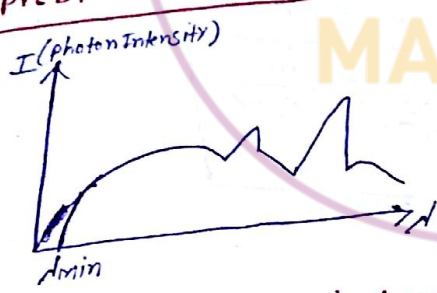
~~Condition~~

Basic requirement of X-ray production

- 1a) → e⁻ producing source
- 1b) → e⁻ Accelerating source (more than 15KV pot.)
- 1c) → Sp. property of target.



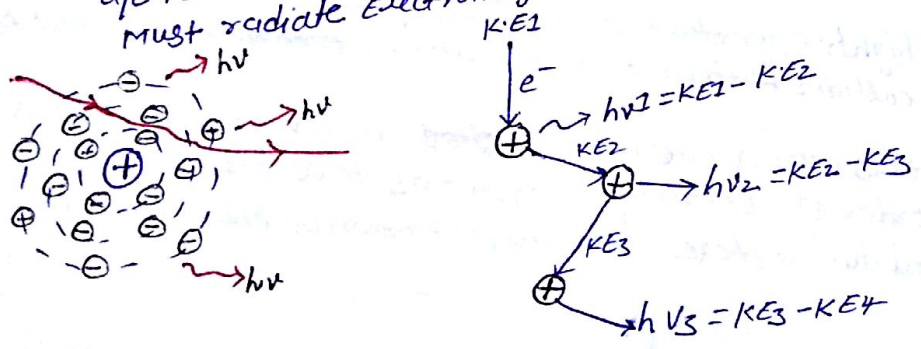
Graph b/w Intensity & Wavelength of Radiation



* There was a wavelength, below which no radiation was obtained, this λ_{min} was named as cutoff wavelength.
 * If we observe the graph, then we see there are two type of variation, one is conti & the other is disconti or, in form of peak.

Production (Explanation) of continuous X-ray / White X-ray.

When highly energetic e^s passed through an atom with very high K.E, they are strongly decelerated in electric field of nucleus & a/c to Maxwell electromagnetic theory every decelerating charge particle must radiate electromagnetic radiation & hence photon are emitted.



In successive collision e^- loses some part of its K.E energy & new photon with lost energy is generated.

** If accelerating voltage across coolidge tube is ' V_0 '
K.E of e^- just before hitting the target

$$\frac{1}{2} m v_e^2 = e V_0$$

$$v_e = \sqrt{\frac{2 e V_0}{m}}$$

** Photon with highest energy is liberated when an e^- lose all its K.E in a single collision & corresponding to this loss only one photon come out.

$$h \nu_{\max} = e V_0 = K.E$$

$$\frac{h c}{\lambda_{\min}} = e V_0$$

$$\text{cutoff wave length} = \lambda_{\min} = \frac{h c}{e V_0}$$

** other photons will have lesser energy so their wavelength will vary from λ_{\min} to ∞ .

$$\Delta \cdot K_{\text{loss}} = E_{\text{x-ray}} = \frac{1}{2} m (v_1^2 - v_2^2)$$

$$\lambda_{\text{x-ray}} = \frac{h c}{E_{\text{x-ray}}} = \frac{h c}{\frac{1}{2} m (v_1^2 - v_2^2)}$$

** Range of $v_2 = 0 - \infty$

$$* v_2 = v_1 \Rightarrow E_{\text{x-ray}} = 0 \Rightarrow \lambda_{\max} = \infty$$

$$* v_2 = 0 \Rightarrow (E_{\text{x-ray}})_{\max} = \frac{1}{2} m v_1^2 = W = e \Delta V$$

$$\lambda_{\min} = \frac{h c}{(E_{\text{x-ray}})_{\max}} = \frac{h c}{\frac{1}{2} m v_1^2} = \frac{h c}{e \Delta V}$$

$$\lambda_{\min} = \frac{12400}{\Delta V (\text{volt})} \text{ \AA}$$

**

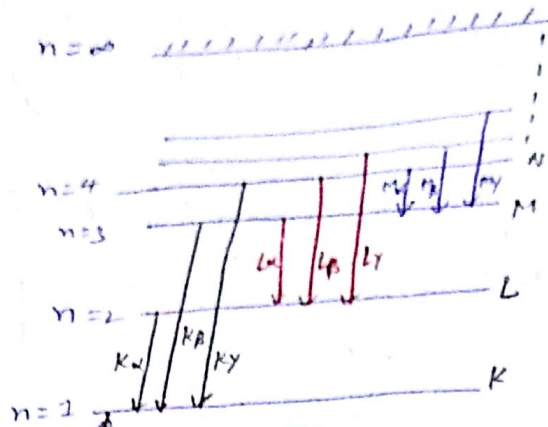
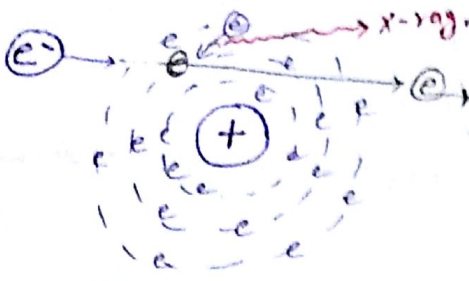
NOTE → Min wavelength of continuous X-ray only depend on anode potential. It is independent from at. no and nature of target.

2016
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Explanation or, production of characteristic X-ray or, explanation of peaks →

When highly energetic e^- enters into an atom, there is also some possibility that it collides with an e^- of the atom & knock out of the shell.

In this way a vacancy is ~~filled~~ created in shell, which is filled by higher order e^- s by making transitions to a lower stage, radiation emitted during these transition is known as characteristic X-ray.



$$V_{K\alpha} < V_{K\beta} < V_{K\gamma}$$

$$\lambda_{K\alpha} > \lambda_{K\beta} > \lambda_{K\gamma}$$

$$V_{K\alpha} > V_{L\alpha} > V_{M\alpha}$$

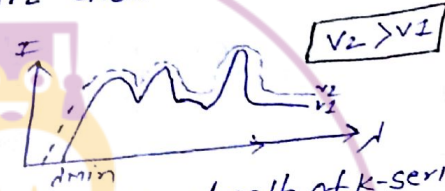
$$\lambda_{K\alpha} < \lambda_{L\alpha} < \lambda_{M\alpha}$$

NOTE → Energy (E_K, E_L, E_M, E_N) are not found from Bohr Model.
 If e^- is knocked out from L-shell.

$$L\alpha, L\beta, L\gamma$$

$$h\nu_{L\alpha} = E_L - E_M$$

$$h\nu_{L\beta} = E_L - E_N$$



- * In a same transition (α, β, γ) Wavelength of α -transition is min.
- * In a same series (K, L, M) Wavelength of α -transition is max.
- * Wavelength of characteristic X-ray doesn't depend on Accelerating voltage, they depend only on target material.

Moseley's Law → Square root of frequency \propto to effective atomic no. of multi e^- system.

$$\sqrt{\nu} \propto Z_{eff}$$

$$\sqrt{\nu} = aZ_{eff}$$

$$Z_{eff} = Z - b$$

↑
At. no.

$$\begin{aligned}
 &K \Rightarrow b \Rightarrow 1 \\
 &L \Rightarrow b \Rightarrow 7.2 \\
 &M \Rightarrow b \Rightarrow 18.2
 \end{aligned}$$

box b
 ↑
 screening const.

Screening const (b) → He explained screening const. as when an e^- makes a transition from a higher shell to lower shell, for eg → K-shell whose a vacancy is created, then the effective nuclear charge ↓ by a factor ' b ' known as screening const.
 In 'K' shell there is only one e^- , which belongs to 's'-orbitals i.e. spherically symmetry so for K-line, effective nuclear charge ~~become~~ become $Z - 1$

* For $K\alpha$ -line

$$\sqrt{\nu_{K\alpha}} = \sqrt{\frac{3RC}{4}} (Z-1)$$

* For $K\beta$ -line

$$\sqrt{\nu_{K\beta}} = \sqrt{\frac{8RC}{9}} (Z-1)$$

* For $K\gamma$ -line

$$\sqrt{\nu_{K\gamma}} = \sqrt{\frac{15RC}{16}} (Z-1)$$

*** → By this observation he concluded that property of material/atom depend on atomic no. not on atomic mass.

$$a = \text{proportional} = \sqrt{Rc \left(\frac{1}{n_1^2} - \frac{1}{n_2^2} \right)}$$

const

NOTE → 'a' depend on transition & 'b' depend on series but both are independent from at. no. of element.

$$\sqrt{V} = a(z-b)$$

$$V = R(z-b)^2 \left(\frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

$$\frac{1}{\lambda} = R(z-b)^2 \left(\frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

$$E = h\nu = Rch = (z-b)^2 \left(\frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

* R = Rydberg const = 10^7 m^{-1}

* $\frac{1}{R} = 912 \text{ \AA}$

* 1 Rydberg = $Rch = 13.6 \text{ eV}$
↑
unit of energy

$$E = 13.6 (z-b)^2 \left(\frac{1}{n_1^2} - \frac{1}{n_2^2} \right) \text{ eV}$$

$$V = 2 \times 10^{25} (z-b)^2 \left(\frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

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A → crystalline solid can cause x-ray to diffract.
 R → Interatomic distance in crystalline solid is of order of 0.2 nm.

Ans → A

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